

**General Relativity Ph236c**  
**Problem Set 3**  
**Due: In class, April 24, 2007**

**Suggested Reading:** Carroll, 8.8; Kolb and Turner, Ch. 8; Dodelson, Ch. 6; Liddle and Lyth, *Cosmological Inflation and Large-Scale Structure*, Ch. 3; Peebles, Ch. 17.

**Preview:** This should be an interesting problem set; in it, you will work out several possibilities, many relevant to current research, for inflation, dark energy, and dark matter. The problems themselves are technically pretty simple, but together they cover quite a bit of territory. Problem 1 is a straightforward exercise in which you will show that a rolling scalar field (i.e., a scalar field in a kinetic-energy-dominated phase) acts like matter with pressure  $p = \rho$ . Problem 2 is a pretty involved problem in which you will work out pretty fully the phenomenological consequences of a particular model for inflation. Problem 3 works through a particularly intriguing quintessence model (i.e., a scalar-field model for negative-pressure dark energy in the Universe today) in which the dark-energy density tracks that of the dominant component (e.g., radiation or matter) of the cosmological energy density. Problem 4 is an order-of-magnitude calculation that shows that magnetic monopoles produced at a GUT phase transition should overwhelm the density of the Universe today (if there were no inflation); this reproduces a calculation that John Preskill was the first to do around 1980. Problem 5 shows that oscillations in an anharmonic scalar-field potential can give rise to exotic equations of state. Problem 6 has you calculate the flux of gamma rays produced by annihilation of WIMPs in the Galactic halo.

1. **A  $w = 1$  equation of state from a rolling scalar field.** Consider a massless scalar field; i.e., a scalar field  $\phi(\vec{x}, t)$  whose potential-energy density is  $V(\phi) = 0$ . Now suppose that this scalar field is initially rolling, so  $\dot{\phi} \neq 0$ , and that the kinetic-energy density associated with this rolling dominates the energy density of the Universe. Show from the stress-energy tensor  $p = \rho$  for this type of matter. Show that this implies that  $\rho \propto a^{-6}$ , where  $a$  is the scale factor, in two ways: (1) by recalling how the energy density of matter with an equation of state  $p = w\rho$  scales with  $a$ ; and (2) by solving the equation of motion for  $\phi$  in an expanding Universe. (This should be a very simple problem.)
2. (From LL 3.7) **Phenomenology of  $\lambda\phi^4$  inflation.** Consider  $V(\phi) = \lambda\phi^4$ , where  $\lambda$  is the self-coupling. Assume that the field rolls toward  $\phi = 0$  from the positive side. Calculate the value of  $\phi$  where each of the slow-roll conditions (i.e.,  $\epsilon \ll 1$  and  $\eta \ll 1$ ) first break down. Do they break down at the same place? Assuming that inflation ends when  $\epsilon = 1$ , calculate the number of  $e$ -foldings of inflation that occur for an initial value  $\phi_i$ . Demonstrate that the slow-roll solutions with  $\phi = \phi_i$  and  $a = a_i$  at  $t = t_i$  are

$$\phi = \phi_i \left[ -\sqrt{\frac{32\lambda M_{\text{Pl}}^2}{6}}(t - t_i) \right],$$

$$a = a_i \exp \left( \frac{\phi_i^2}{8M_{\text{Pl}}^2} \left\{ 1 - \exp \left[ -\sqrt{\frac{64\lambda M_{\text{Pl}}^2}{3}}(t - t_i) \right] \right\} \right).$$

Use the solution for  $\phi$  to calculate the time that inflation ends. Demonstrate that the number of  $e$ -foldings calculated using the solution for  $a$  is the same as that which you calculated above. Expand the solution for  $a$  at small  $t - t_i$  to demonstrate that the inflation is approximately exponential in the initial stage. Calculate the time constant  $\kappa$  [from  $a \sim \exp(\kappa t)$ ] and demonstrate that it equals the (slow-roll) Hubble parameter during inflation.

3. **Tracker field.** Consider a scalar field that rolls down a potential-energy density  $V(\phi) = V_0 e^{-\phi/\phi_0}$ . Now suppose that the energy density of the Universe is dominated by ordinary non-relativistic matter (so  $a \propto t^{2/3}$ ), and that the energy density of the rolling scalar field is negligible compared with the non-relativistic matter. Show that there is a solution to the scalar-field equation of motion such that the energy density  $\rho_\phi = (1/2)\dot{\phi}^2 + V(\phi)$  of the scalar field scales as  $\rho_\phi \propto a^{-3}$ , the same as the ordinary matter. Does the same thing happen if the energy density of the Universe is dominated by relativistic matter? This is the basis for the “tracker-field” solutions that have been discussed in the literature recently.
4. **The monopole problem.** Calculate the relic density of magnetic monopoles, assuming that there is one GUT-mass ( $\sim 10^{15}$  GeV) monopole produced per Hubble volume at the GUT phase transition ( $T \sim 10^{15}$  GeV). You should get an unreasonably large number. There is a bound  $\Omega_{\text{monopole}} \lesssim 10^{-6}$  (the Parker bound) to the relic density of magnetic monopoles in the Universe today. Calculate the number of  $e$ -folds of inflation after the GUT transition required to solve the monopole problem.
5. **Anharmonic scalar-field oscillations.** In class we argued that if we have a real scalar field  $\phi$  with a quadratic potential  $V(\phi) = (1/2)m^2 H^2$ , and if  $m \gtrsim H$  (implying that the oscillation frequency is large than the expansion rate), then coherent oscillations of the scalar field imply that the pressure  $p = 0$  when averaged over an oscillation cycle and thus that the energy density  $\rho \propto a^{-3}$ . Now consider oscillations in a potential  $V(\phi) = c|\phi|^n$ , where  $c$  is a constant. Show that coherent oscillations in such a potential give rise to an energy density that decays as  $\rho \propto a^{-\alpha}$ , and determine  $\alpha$ . Of course, you should recover  $\alpha = 3$  for  $n = 2$ . What value of  $n$  is required to produce  $\alpha = 4$  (i.e., radiation)? Can you think of a physical argument that justifies your result? Likewise, is there a value of  $n$  that produces  $\alpha = 0$ ? Can you explain this result in physical terms?
6. **WIMPs in the Galactic halo.** This problem has nothing to do with inflation, but goes back to weakly-interacting massive particle (WIMP) dark matter. Suppose the Galactic halo has a density distribution,

$$\rho(r) = \rho_0 \frac{r_0^2 + a^2}{r^2 + a^2},$$

where  $\rho_0 = \rho(r_0)$  is the local halo density,  $r_0 \simeq 8.5$  kpc is our distance from the Galactic center, and  $a \simeq 4$  kpc is the core radius. Suppose also that the halo is made out of

WIMPs  $\chi$  with masses  $m_\chi = 10\text{s} - 1000\text{s GeV}$  which have a cross section for annihilation to two photons  $\sigma_{\chi\chi \rightarrow \gamma\gamma}$ . (a) If  $\psi$  is the angle some given line of sight makes with the Galactic center (i.e.,  $\psi = 0$  is toward the Galactic center and  $\psi = \pi$  is away from the Galactic center), calculate the flux of gamma rays observed along the line of sight  $\psi$ . (b) What is the energy of these photons. (c) The WIMP  $\chi$  may annihilate to photons via a diagram such as that below, where the particles in the loops are quarks  $q$  and squarks  $\tilde{q}$ , and the squark mass is comparable to the WIMP mass. Give an order-of-magnitude estimate of this annihilation cross section and estimate the flux of photons from WIMP annihilation.