Quantum Field Theory (171.701), Fall 2022

Problem Set 2

Due: 7 November 2022

1. Consider the Landau-Ginsberg free energy for a complex scalar ψ ,

$$F = \int dx \left(\alpha_2(T) |\psi|^2 + \alpha_4 |\psi|^4 - \gamma \left| \frac{d\psi}{dx} \right|^2 + \kappa \left| \frac{d^2\psi}{dx^2} \right|^2 \right), \qquad (1)$$

with $\gamma, \alpha_4, \kappa > 0$. Consider an ansatz in which a single Fourier mode $\psi_k = A_k e^{ikx}$ with $k = \pm k_0$ is nonvanishing. What value of k_0 minimizes the free energy? Show that the system undergoes a phase transition to a spatially modulated phase when $\alpha_2 = \gamma^2/4\kappa$. What symmetries are broken in the ordered phase?

2. Consider the free energy for a complex scalar field ψ coupled to a gauge field A_i in d dimensions,

$$F[\psi, A_i] = \int d^d x \, \left[\frac{1}{4} F_{ij} F^{ij} + |\partial_i \psi - ieA_i \psi|^2 + \mu^2 |\psi|^2 \right], \qquad (2)$$

where $F_{ij} = \partial_i A_j - \partial_j A_i$. (As an aside: if we add a quartic term $g|\psi|^4$, then, in d = 3, this is the original Ginzburg-Landau free energy for a superconductor.) What is the critical dimension d_c , such that the coupling between the scalar and gauge field is relevant for $d < d_c$ and irrelevant for $d > d_c$? Speculate on what might happen at $d = d_c$.

3. The O(N) mode consists of N scalar fields $\phi = (\phi_1, \dots, \phi_N)$, with a free energy,

$$F_{1}[\phi] = \int d^{d}x \, \left[\frac{1}{2} \left(\nabla\phi\right)^{2} + \frac{1}{2}\mu^{2} \left(\phi \cdot \phi\right) + g \left(\phi \cdot \phi\right)^{2}\right], \quad (3)$$

which is invariant under O(N) rotations. In $d = 4 - \epsilon$ dimensions, the beta functions become,

$$\frac{d\mu^2}{ds} = 2\mu^2 + \frac{N+2}{2\pi^2} \frac{\Lambda^4}{\Lambda^2 + \mu^2} \tilde{g},
\frac{d\tilde{g}}{ds} = \epsilon \tilde{g} - \frac{N+8}{2\pi^2} \frac{\Lambda^4}{(\Lambda^2 + \mu^2)^2} \tilde{g}^2,$$
(4)

where $\tilde{g} = \Lambda^{-\epsilon} g$ is the dimensionless coupling constant. What is the critical exponent ν at the Wilson-Fisher fixed point. Assuming that $\eta \sim \mathcal{O}(\epsilon^2)$, determine the critical exponents α , β , γ , and δ to leading order in ϵ .

4. The free energy of the Sine-Gordon model in d = 2 dimensions, with UV cutoff Λ , is given by

$$F[\phi] = \int d^2x \left[\frac{1}{2} (\nabla \phi)^2 - \lambda_0 \cos(\beta_0 \phi) \right], \qquad (5)$$

What is the "naive" dimension of β_0 and λ_0 . Decompose the modes of the field as $\phi_{\mathbf{k}} = \phi_{\mathbf{k}}^- + \phi_{\mathbf{k}}^+$, where ϕ_k^+ is nonzero for $Lambda/\zeta < k < \Lambda$, and ϕ_k^- nonzero for other k. Let $\phi^-(\mathbf{x})$ and $\phi^+(\mathbf{x})$ be the inverse Fourier transforms. Show that after integrating out $\phi_{\mathbf{k}}^+$, the free energy for $\phi^$ becomes, to leading order in λ_0 ,

$$F'[\phi^{-}] = \int d^2x \, \left[\frac{1}{2} (\nabla \phi^{-})^2 - \lambda_0 \left\langle \cos \beta_0 (\phi^{-} + \phi^{+}) \right\rangle_+ \right], \tag{6}$$

where you should define the meaning of $\langle \, \rangle_+.$ Evaluate this expectation value to show that

$$F'[\phi^{-}] = \int d^2x \, \left[\frac{1}{2} (\nabla \phi^{-})^2 - \lambda_0 \zeta^{-\beta_0^2/4\pi} \cos(\beta_0 \phi^{-}) \right]. \tag{7}$$

Hence, show that the $\cos(\beta\phi)$ potential is relevant when $\beta_0^2 < 8\pi$ and irrelevant when $\beta_0^2 > 8\pi$. [Hint: Write $\cos(\phi^- + \phi^+) = \frac{1}{2}(e^{i\phi^-}e^{i\phi^+} + e^{-i\phi^-}e^{-i\phi^+})$ and use Wick's identity,

$$\left\langle e^{B_a\phi_a}\right\rangle = e^{\frac{1}{2}B_a\left\langle\phi_a\phi_b\right\rangle B_b},\tag{8}$$

for any constant B_a . You will also need the position-space correlation function,

$$\left\langle \phi^{+}(\mathbf{x})\phi^{+}(\mathbf{y})\right\rangle_{+} = \int_{\Lambda/\zeta}^{\Lambda} \frac{d^{2}k}{(2\pi)^{2}} \frac{e^{-i\mathbf{k}\cdot(\mathbf{x}-\mathbf{y})}}{k^{2}}.$$
 (9)

You should make use of this in the limit $\mathbf{x} \to \mathbf{y}$.]

5. Consider a field φ_i in a representation R_1 of some gropu and a field χ_I in a representation R_2 . Their product $\phi_i \chi_I$ is then in the direct product representation $R_1 \otimes R_2$, with generator matrices given by

$$(T^{a}_{R_{1}\otimes R_{2}})_{iI,jJ} = (T^{a}_{R_{1}})_{ij} \,\delta_{IJ} + \delta_{ij} \,(T^{a}_{R_{2}})_{IJ} \,. \tag{10}$$

(a) Prove the distribution rule for the covariant derivative,

$$[D_{\mu}(\varphi\chi)]_{iI} = (D_{\mu}\varphi)_i \chi_I + \varphi_i (D_{\mu}\chi)_I.$$
⁽¹¹⁾

(b) Consider a field φ_i in a representation and its Hermitian conjugate φ_i^{\dagger} in the complex-conjugate representation. Show that

$$\partial_{\mu}(\varphi^{\dagger i}\varphi_{i}) = (D_{\mu}\varphi^{\dagger})^{i}\varphi_{i} + \varphi^{\dagger i}(D_{\mu}\varphi).$$
(12)

Explain why this is a special case of equation (11), and also why it makes sense.